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Aiming factor to flatten the flux distribution on cylindrical receivers

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Abstract

High incident flux gradients and hot spots lead to extreme thermal stresses that may damage and reduce the lifetime of central receivers. An aiming strategy based on a single parameter, k , named aiming factor, is developed to generate symmetric flux maps about the receiver equator. By means of this k factor, ranging between 3 (generally equivalent to equatorial aiming) and 0 (alternatively aiming to top and bottom borders), the solar flux incident on the receiver and the spillage losses can be controlled. For each sector in a heliostat field, the aiming factor values causing the flattest symmetric flux maps, k_{flat} , are deterministically found with a sweep and mesh shifting procedure. Results for Dunhuang solar power tower plant show that k_{flat} is fairly constant throughout the year, except near sunrise and sunset in east and west sectors, respectively.

Keywords: Solar power tower, Heliostat spillage, Symmetric flux map, Flat concentration profile, FluxSPT tool

1 Nomenclature

AC	Cell area in the receiver mesh [m ²]
AFD	Allowable flux density [W/m ²]
AM	Mirror area of the heliostat [m ²]

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BR_k	Beam radius based on k [m]
C	Concentration ratio of flux density [-]=[suns]
D	Receiver diameter [m]
DNI	Direct normal irradiation [W/m ²]
F	Flux density [W/m ²]
f	Loss factor [-]
H	Receiver height [m]
k	Aiming factor [-]
\mathbf{n}	Heliostat normal vector
$RMSD$	Root Mean Square Deviation
\mathbf{s}	Sun vector
SR	Slant range [m]
\mathbf{t}	Target vector
X, Y, Z	Cartesian coordinate axes

2 *Greek symbols*

ϵ	Elevation angle [rad]
η	Efficiency [-]
ω	Incidence angle [rad]
σ	Gaussian error [mrad]

3 *Subscripts*

e	Effective
h	Heliostat
i, j, p	Nodes in receiver mesh
max	Maximum
si	Image plane system of coordinates
slp	Slope
sp	Spillage
st	Target system of coordinates
sun	Sunshape
trk	Tracking

4 *Acronyms*

DAPS	Dynamic Aim Processing System
FluxSPT	Flux Solar Power Tower software tool
MCRT	Monte Carlo Ray Tracing
PID	Proportional Integral Derivative controller
PS10	Planta Solar 10 MW
PSA	Plataforma Solar de Almería
SAPS	Static Aim Processing System
SPT	Solar Power Tower
UNIZAR	Universidad de Zaragoza tool

5 1. Introduction

6 In Solar Power Tower (SPT) plants, direct radiation is concentrated by
7 thousands of heliostats onto a tower-mounted central receiver, where a cir-
8 culating working fluid is heated to eventually produce electricity [1]. Single
9 equatorial aiming leads to the highest receiver interception, but also to un-
10 acceptable peak fluxes that have to be avoided or minimized. For successful
11 operation of these plants, the heliostat field aiming strategy must protect the
12 receiver from damage (thermal stress cracking, hot spots and corrosion), at
13 the same time that the thermal output is maximized.

14 As long as SPT technology is being commercially deployed worldwide,
15 the development of efficient aiming strategies is receiving great attention
16 from the academic community in recent years. Metaheuristic methods, such
17 as Tabu search and Genetic Algorithm, were applied to uniform the flux
18 distribution in flat plate [2, 3] and cavity receivers [4, 5]. By means of an ant
19 colony metaheuristic, the output of concentrated photovoltaic receivers [6]
20 and molten salt receivers [7] was optimized. Astolfi et al. [8] proposed several
21 optimization approaches to reduce the peak flux in a cylindrical receiver. A
22 teaching-learning-based optimization was investigated to homogenize the flux
23 on a flat receiver [9], and a closed-loop PID control was virtually implemented
24 on Gemasolar plant [10]. Likewise, binary integer linear programming was
25 virtually implemented on PS10 plant to homogenize the flux distribution on
26 the flat plate receiver [11].

27 In Solar Two plant, two software systems controlled the field aiming:
28 SAPS (static aim processing system) and DAPS (dynamic aim processing
29 system). Using the first system, the target point of each heliostat was ad-
30 justed every 10 minutes because of sun movement, whereas heliostats caus-

31 ing overheating were removed from tracking with DAPS [12]. In PSA tower
32 plant, named CESA-1, a knowledge-based heuristic selected the optimal aim
33 points to control the air temperature in the volumetric receiver [13]; and,
34 a distributed optimization algorithm was recently performed to uniform the
35 flux distribution [14].

36 In order to procure symmetric flux maps, Vant-Hull suggested an aiming
37 strategy where each row of heliostats is alternatively aimed to the top and
38 to the bottom half of the receiver [15]. This strategy was further developed
39 by the authors in Ref. [16], where it was succinctly introduced the concept
40 of aiming factor.

41 The aiming factor approach [16], that has been utilized in recent refer-
42 ences by these [17, 18] and other [10, 19] authors, relies on a single parameter
43 to aim complete fields of heliostats. Because of the close relation between
44 k aiming factor and spillage losses, the preliminary objective of the present
45 work is to provide a thorough analysis on the influence of k factor on flux
46 profiles and, ultimately, on spillage losses.

47 Recent research evidences that uniform flux distribution in receiver tubes
48 plays a key role in reducing the peak thermal stress [20], which is the most
49 limiting factor in the most irradiated receiver panels [17]. In this respect, the
50 main purpose of the work is to formerly find the flattest possible flux distri-
51 butions, objective for which the symmetric aiming methodology is adopted.

52 The manuscript is structured as follows. First, for a single heliostat it
53 is defined the aiming factor in order to estimate the size of the beam and
54 consequently point the heliostat. Dunhuang SPT plant, case study used
55 throughout this work, is introduced in Section 3. Section 4 analyzes the
56 effect of the aiming factor approach on the receiver spillage. For the whole
57 field of heliostats, symmetric aiming is applied in Section 5. And Section 6
58 explores the aiming to achieve uniform flux profiles. At the end, the Appendix
59 describes the flux map shifting procedure.

60 The flux distributions by heliostat fields have been computed with the
61 convolution-projection method reported in Ref. [16] and experimentally vali-
62 dated in Refs. [16, 21]. The resulting in-house software FluxSPT is available
63 for free download from the link in Ref. [22]. FluxSPT currently contains three
64 existing SPT plants of moderate, medium and large size, namely: Dunhuang,
65 Gemasolar and Crescent Dunes.

66 **2. Aiming factor**

67 The proposed aiming strategy is based on a single parameter, k , named
 68 aiming factor. For a single heliostat, this Section describes: first, the estima-
 69 tion of the beam size as a function of k factor; and, then, the determination of
 70 the aim point. This aiming strategy is implemented in the FluxSPT software
 71 tool [22].

72 *2.1. Beam radius*

73 The present aiming strategy relies on a proper estimation of the size of
 74 the beam incident on the receiver. In principle, a particular size cannot be
 75 defined for the beam reflected by a heliostat. In the following, it is introduced
 76 a calculation procedure on the basis of k aiming factor.

77 The flux density distribution on the image plane produced by a focus-
 78 ing heliostat follows an essentially circular Gaussian distribution with effec-
 79 tive standard deviation σ_e . This evidence, supported by measurements and
 80 MCRT simulations, is inherently included in well-known convolution models
 81 as: UNIZAR [23], DELSOL [24] and HFLCAL [25].

82 Herein, it has been adopted the analytic function on the image plane
 83 by UNIZAR, where the effective standard deviation (σ_e) results from the
 84 convolution of sunshape (σ_{sun}), mirror slope (σ_{slp}), and tracking (σ_{trk}) errors.
 85 In the effective error equation, ω_h stands for the incidence angle on the
 86 heliostat.

$$\sigma_e = \sqrt{\sigma_{sun}^2 + 2(1 + \cos \omega_h) \sigma_{slp}^2 + \sigma_{trk}^2} \quad (1)$$

87 By analogy with a circular normal distribution, 68%, 95% and 99.7% of
 88 the total flux is within the cone of aperture angle from heliostat center equal
 89 to σ_e , $2 \cdot \sigma_e$ and $3 \cdot \sigma_e$, respectively. Therefore, it can be defined a factor k ,
 90 generally ranging between 0 and 3, that gives an idea on the energy that lies
 91 within $k \cdot \sigma_e$, in line with the 68-95-99.7 rule for normal distribution.

92 The radius of the beam (BR) on the image plane (si), that normal to the
 93 main reflected ray (or \mathbf{t} target vector), is derived from the cone geometry.
 94 For a given SR slant range (distance from the heliostat to the receiver), the
 95 beam radius is function of the k factor, as declared in Eq. 2. Obviously, the
 96 higher the k value, the larger the estimated beam circle is.

$$BR_k^{si} = SR \cdot \tan(k \cdot \sigma_e) \simeq SR \cdot k \cdot \sigma_e \quad (k \cdot \sigma_e \ll 1rad) \quad (2)$$

97 Fig. 1 displays the flux distribution produced by a heliostat in both a 3D
 98 view (left) and a 2D view (center), as well as the vertical profile through the

99 target point (right). Red circles outline the beam circumferences for three k
100 values: 1, 2 and 3. Instead of flux density, F , in W/m^2 , the distribution is
101 quantified in concentration ratio of flux density, C , taking advantage of its
102 independence of instantaneous direct normal irradiation, DNI . The concen-
103 tration ratio of flux density is dimensionless, as derived from Eq. 3, while
104 also can be expressed as the number of instantaneous suns focusing on the
105 receiver.

$$F = C \cdot DNI \quad (3)$$

106 The flux distribution on the image plane is shown in Fig. 1a. The red
107 circles are the beam boundaries for $k = 3, 2, 1$, whose radii are estimated
108 with previous Eq. 2.

109 On the surfaces of the receiver, the flux distribution is projected from the
110 image plane, following the procedure of the optical model, that is detailed
111 in Ref. [16]. Similarly, the beam circumference is projected on the receiver
112 panels, which is equivalent to the intersection of the cone of light with the
113 receiver. The flux distribution on the half visible side of the receiver is shown
114 in Fig. 1b.

115 For aiming strategy purposes, the beam size on the vertical of the receiver
116 is only needed. The radius of the beam in vertical direction as a function of
117 k factor, BR_k , can be easily calculated from that on the image plane BR_k^{si} .
118 Known the elevation angle of the target vector (ϵ_t), Eq. 4 estimates the beam
119 radius depending on k factor.

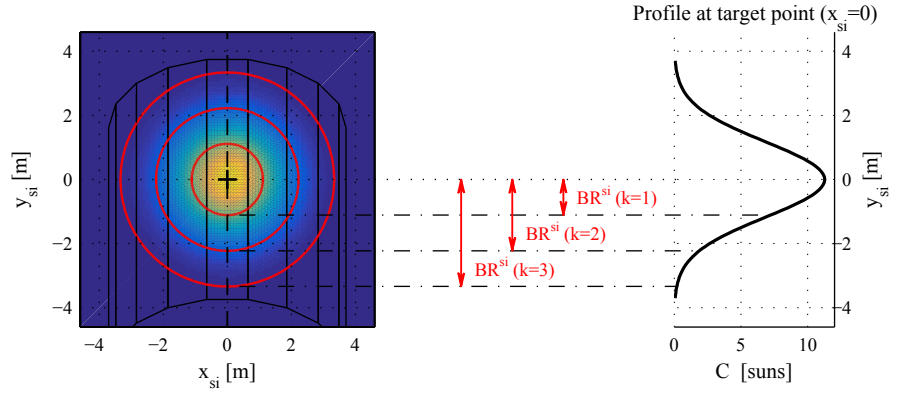
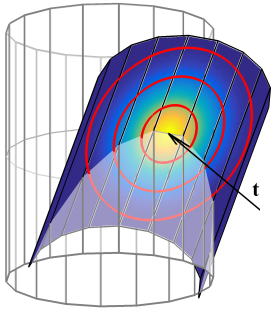
$$BR_k = \frac{BR_k^{si}}{\cos \epsilon_t} = \frac{SR \cdot k \cdot \sigma_e}{\cos \epsilon_t} \quad (4)$$

120 2.2. Targeting

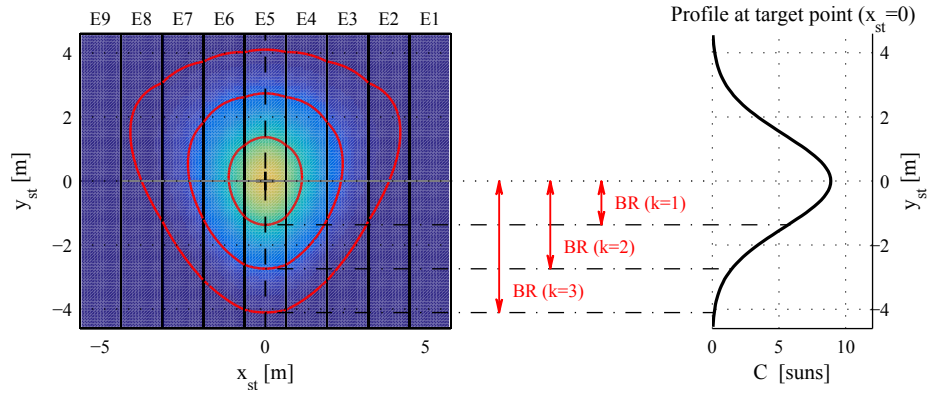
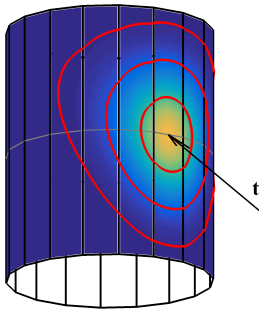
121 Once the beam radius has been estimated as a function of the aiming
122 factor, the heliostat is targeted in such a way that the beam is tangent
123 to either the lower or the upper edge of the receiver. In other words, the
124 heliostat is pointed at a vertical distance BR_k from one of the edges.

125 For the sample heliostat used in the preceding Subsection, Fig. 1c il-
126 lustrates the concentration map that results from applying $k = 2$. In this
127 example, the beam is tangent to the bottom receiver edge.

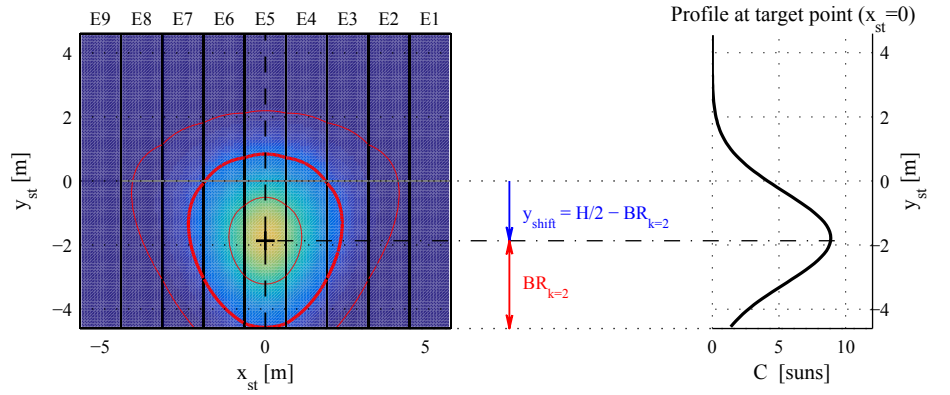
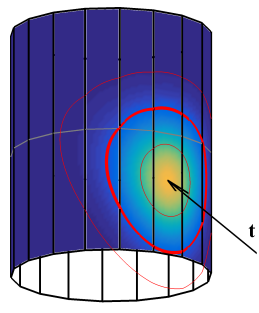
128 This targeting procedure entails the adjustment of the target point along
129 the vertical direction on the receiver surface. To minimize lateral spillage
130 losses in cylindrical receivers, the azimuth of the aim point is that of the
131 pointing heliostat; i.e. the target vector is coplanar with the vertical axis of



(a) On image plane.



(b) On receiver surface.



(c) Targeting for $k = 2$.

Image at summer solstice noon
by Dunhuang heliostat in 7th row, east azimuth.

Figure 1: Flux distribution and beam boundaries (in red) for several aiming factors $k = 3, 2, 1$.

132 the receiver/tower. Thus, vertical shifting is only considered, in contrast to
 133 horizontal shift which is not supported.

134 Whether the beam diameter, given a k factor, is greater than the receiver
 135 height (H), equatorial aiming is kept; on the contrary, spillage losses would
 136 unwillingly increase. The vertical shifting of the aim point respect to the
 137 equatorial target point (y_{shift}) is accordingly summarized in the following
 138 Equation. The minus sign applies for aiming tangent to the lower edge,
 139 otherwise the shift is upwards.

$$y_{shift} = \begin{cases} \pm \left(\frac{H}{2} - BR_k \right) & \text{if } BR_k < H/2 \\ 0 & \text{if } BR_k \geq H/2 \end{cases} \quad (5)$$

140 Since the beam size decreases with k aiming factor, then the lower the
 141 k factor, the higher the vertical shift is. This is why parameter k has been
 142 named aiming factor.

143 3. Case study

144 The methodology presented in this paper is valid for any SPT plant with
 145 central cylindrical receiver. As case study to illustrate the procedure and to
 146 show results, it is considered Dunhuang SPT plant. This is a 10 MW_e plant
 147 located in China at 40.08° north latitude.

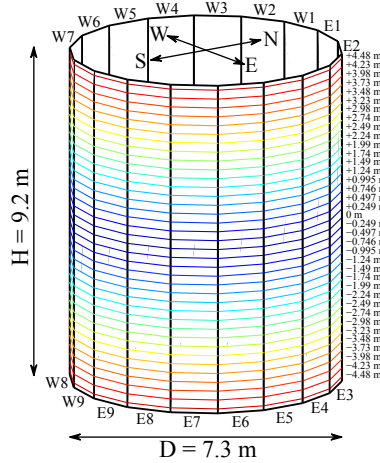


Figure 2: Dunhuang receiver geometry and aim levels under consideration.

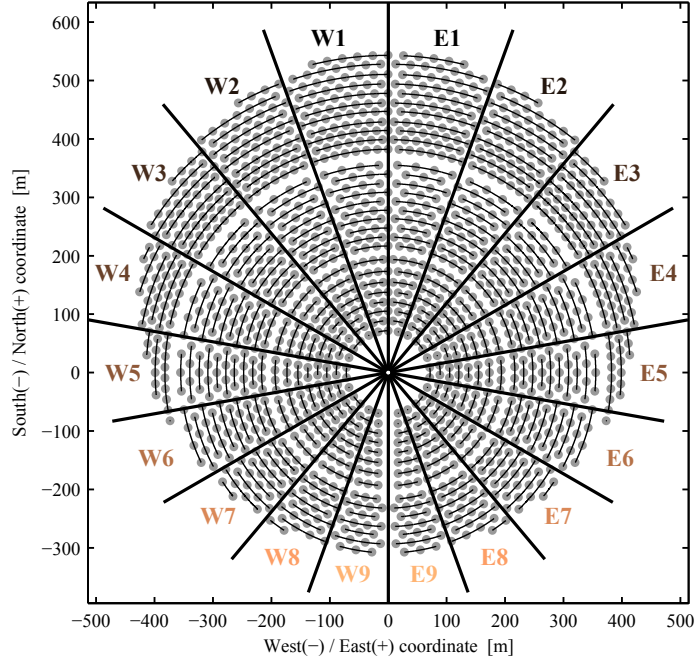


Figure 3: Heliostat field layout of Dunhuang with divisions by sector.

148 The receiver, 9.2 m in height and 7.3 m in diameter, comprises 18 panels,
 149 and stands at the top of the tower, being 121 m its optical height. Fig. 2 de-
 150 picts the receiver, where the panels are labeled according to their orientation
 151 (East or West), and numbered from the north. Aim levels to be considered
 152 are also marked with colored lines in the Figure.

153 The field consists of 1525 square heliostats of 115 m² mirror surface each.
 154 Fig. 3 depicts the heliostat field layout, and its division in sectors according
 155 to their target panels.

156 For the evaluation of the effective error (Eq. 1), a reliable value of sun-
 157 shape standard deviation equal to 2.09 mrad has been taken [26, 27]. A
 158 common error of 2.6 mrad is handled for the mirror slope error, while the
 159 tracking error is neglected since it does not produce any actual modification
 160 on the flux distribution by the heliostat.

161 4. Effect on spillage

162 This Section analyzes how the above described aiming factor approach
 163 affects the receiver spillage. Actually, k parameter not only determines the

164 target point of each heliostat in a field, but also gives some idea of the mag-
 165 nitude of the spillage losses, or its complimentary, the receiver interception.

166 Spillage losses are quantified by the spillage factor (f_{sp}), also known as
 167 intercept factor. This loss factor is the fraction of direct solar flux reflected
 168 by the heliostat that is actually intercepted by the receiver. In terms of the
 169 concentration ratio in each node of the receiver, $C_{i,j,p}$ in receiver mesh node
 170 coordinates, the spillage factor can be computed with the following Equation
 171 [16], where: AC stands for the area of the cells in the receiver mesh, and AM
 172 is the mirror area of the heliostat.

$$f_{sp} = \frac{\sum_{i,j,p} C_{i,j,p} \cdot AC}{\cos \omega_h \cdot AM} \quad (6)$$

173 In the following examples, the equinox day at solar noon time is consid-
 174 ered. Fig. 4 shows the spillage factors of the heliostats in Dunhuang field at
 175 equinox noon for equatorial aiming, and taking aiming factors equal to 2, 1
 176 and 0. In this example, every single heliostat is targeted so that the beam
 177 circumference, whose radius is computed with Eq. 4, is tangent to the bot-
 178 tom edge of the receiver; what is termed downwards aiming in the following.
 179 Half of the field is displayed because of symmetric situation at solar noon.

180 Evidently, as the aiming factor decreases the spillage losses increase. As
 181 shown in the Figure, receiver interception for $k = 2$ is only very slightly
 182 lower than with equatorial aiming. Indeed, aiming with $k = 3$ is equivalent
 183 to equatorial aiming for Dunhuang SPT plant.

184 From the basics of the proposed aiming strategy, a given aiming factor
 185 would produce the same spillage losses for all the heliostats in a field. This
 186 fact is noticeable for small aiming factors, but not fully for high k values;
 187 where the spillage colormap is not completely uniform in Fig. 4. Beam spots
 188 greater than the receiver size in distant heliostats justifies spillage factors
 189 smaller than 1, even for equatorial aiming.

190 A more detailed analysis on the influence of the aiming factor on the
 191 spillage losses is provided in Fig. 5 for selected heliostats, specifically those
 192 circled in black in previous Fig. 4. The Figure depicts both the cases of
 193 downwards and upwards aiming, in respectively solid and dashed lines.

194 For twelve heliostats in row 11, Fig. 5a represents k vs. f_{sp} , where is
 195 again evidenced that the lower the k factor, the higher the spillage losses
 196 are. At equinox noon in Dunhuang plant, given a k value, the spillage losses
 197 are moderately constant, regardless of the heliostat azimuthal position.

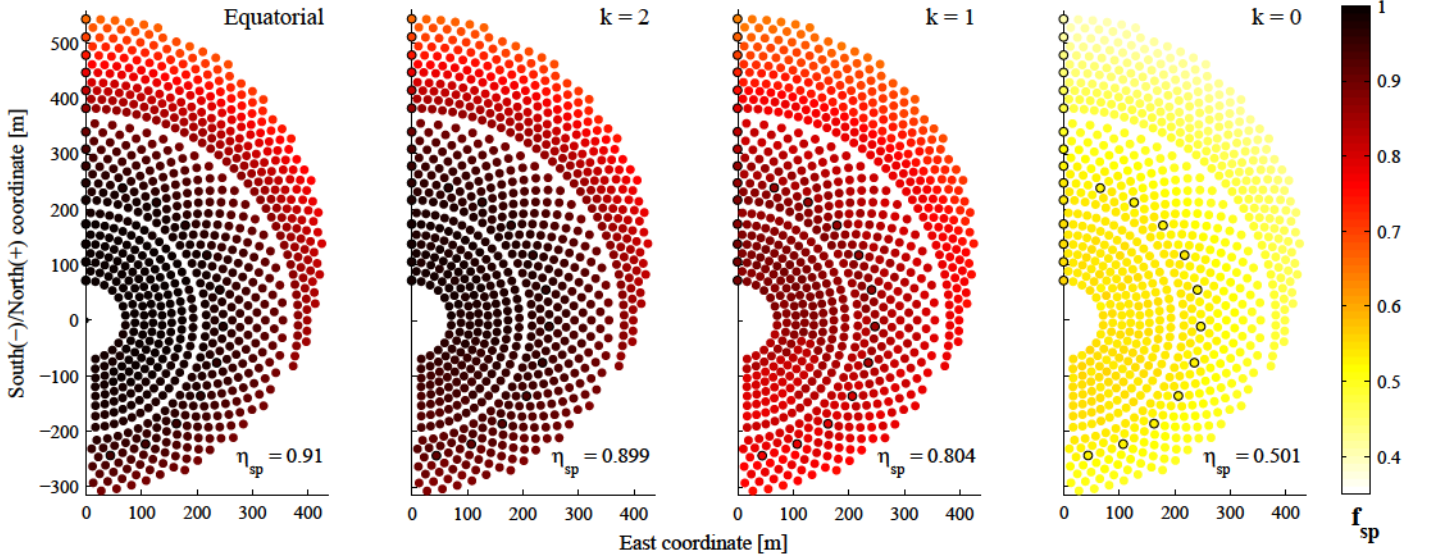
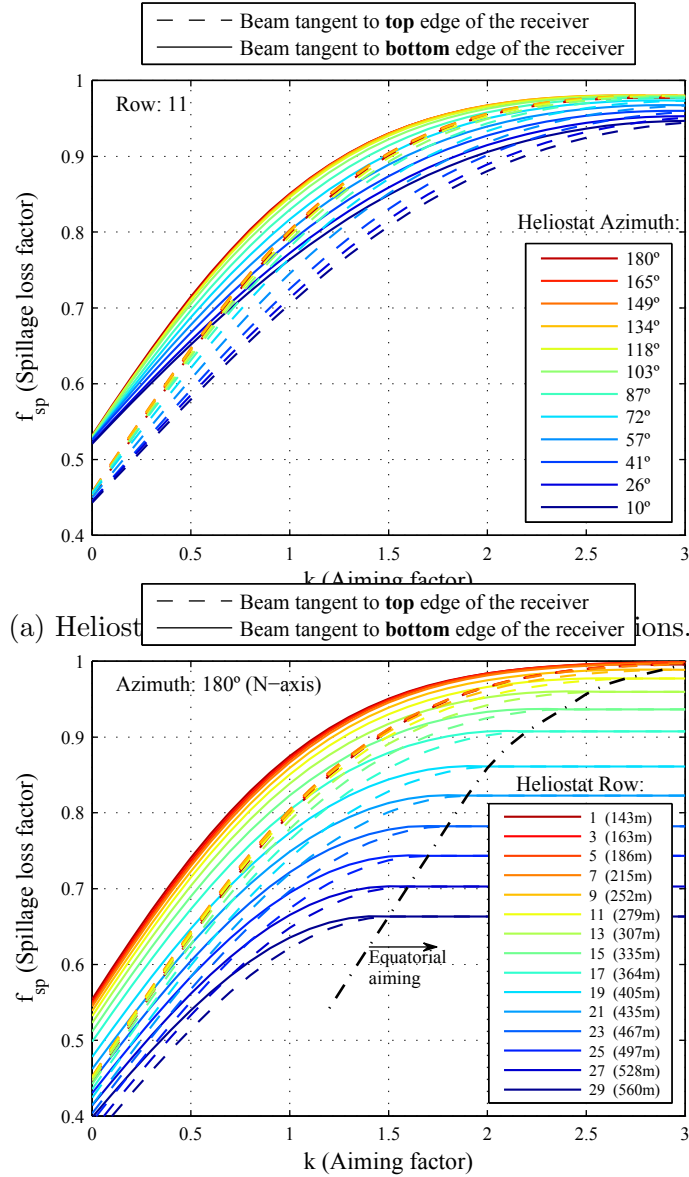


Figure 4: Heliostat spillage loss factors in Dunhuang field (east half) at equinox noon for equatorial aiming, and aiming factors: 2, 1 and 0. Beams tangent to bottom receiver edge (downwards aiming).

198 Fig. 5b similarly represents the cases of fifteen heliostats in the north
 199 axis. Given a k value, now the spillage losses are diverse depending on the
 200 heliostat distance to the tower. From mid-distance rows to last ones, it is
 201 noticed that the farther the heliostat, the higher the spillage loss is. This
 202 amplitude is remarkable for high aiming factors, because of the larger beam
 203 sizes for last rows compared to the receiver size. As can be seen in the Figure,
 204 the spillage loss is maximum (flat top) from a k value onward; such situation
 205 points out that equatorial aiming is taking place. Dash-dot line in Fig. 5b
 206 demarcates the equatorial aiming area.

207 From Fig. 5 it is concluded that the progression of spillage losses as a
 208 function of the aiming factor follows a rather logarithmic profile. Whether a
 209 discrete sweep of k factors is performed (e.g. in Subsection 6.2 to determine
 210 flat flux profiles), a logarithmically spaced scan of k would roughly lead to
 211 constant steps in the spillage losses; as well as in the aim point shifting.

212 By comparing for a given heliostat and k factor, both downwards and
 213 upwards aiming in Fig. 5, respectively in solid and dashed lines, downwards
 214 aiming results in higher interception efficiency; except for equatorial aiming
 215 or high k values, where both tend to match. Even though the beam is placed



(b) Heliostats with azimuth 180 (north axis) in odd rows.

Figure 5: Spillage loss factor of selected heliostats in Dunhuang plant at equinox noon depending on aiming factor. The selected heliostats are circled in black in Fig. 4. Solid and dashed lines correspond to respectively downwards and upwards aiming.

216 tangent to the receiver border (bottom or top) in both cases, since the helio-
217 stat reflected beam follows an ascending direction, receiver panels adjacent
218 to the target one are on average reached above the aim level. Hence, receiver
219 interception is smaller when upwards aiming is set. This is also evident in
220 the spillage efficiency of the whole heliostat field (η_{sp}), from inspection of
221 upwards and downwards outputs in Table 1.

222 5. Symmetric aiming

223 5.1. Methodology

224 From a whole field of heliostats, symmetric flux maps about the receiver
225 equator can be achieved by alternatively aiming each row of heliostats up-
226 wards and downwards. Such approach was formerly proposed by Vant-Hull
227 in Ref. [15].

228 In the present study, even rows of heliostats aim tangentially to the bot-
229 tom receiver edge (downwards), while odd rows aim to the top edge (up-
230 wards). This symmetric aiming strategy is graphically depicted in Fig. 6,
231 where not only the resulting flux map, but also the aiming points are sym-
232 metric about the receiver equator. The summation of the flux maps from
233 upwards (odd rows) and downwards (even rows) aiming leads to a nearly
234 symmetric flux map.

235 The symmetric aiming procedure illustrated in Fig. 6 is provided for the
236 Dunhuang case study, considering equinox day at 8:00 solar time. Aiming
237 points and flux map have been computed for an aiming factor equal to 1.5 in
238 all the heliostats of the field. Parallel projection between neighbor heliostats
239 has been utilized in the calculation of shading and blocking losses.

240 The determination of aim points from a given k factor, following the
241 procedure described in Subsection 2.2, results in a continuous space of aim
242 levels. To simplify the control of the heliostat field, it is generally defined a
243 discrete number of aim levels on the receiver. Then, each and every heliostat
244 is pointed to the aim level whose distance to the equator is immediately
245 below the vertical shift computed with continuous Eq. 5; which can be called
246 shifting quantization. It becomes clear that the higher the number of aim
247 levels, the more possibilities of aiming exists and flatter flux profiles are
248 achievable.

249 For the Dunhuang case study, it has been considered 37 aim levels; an odd
250 number is required to attach the aim level in the middle to the equatorial

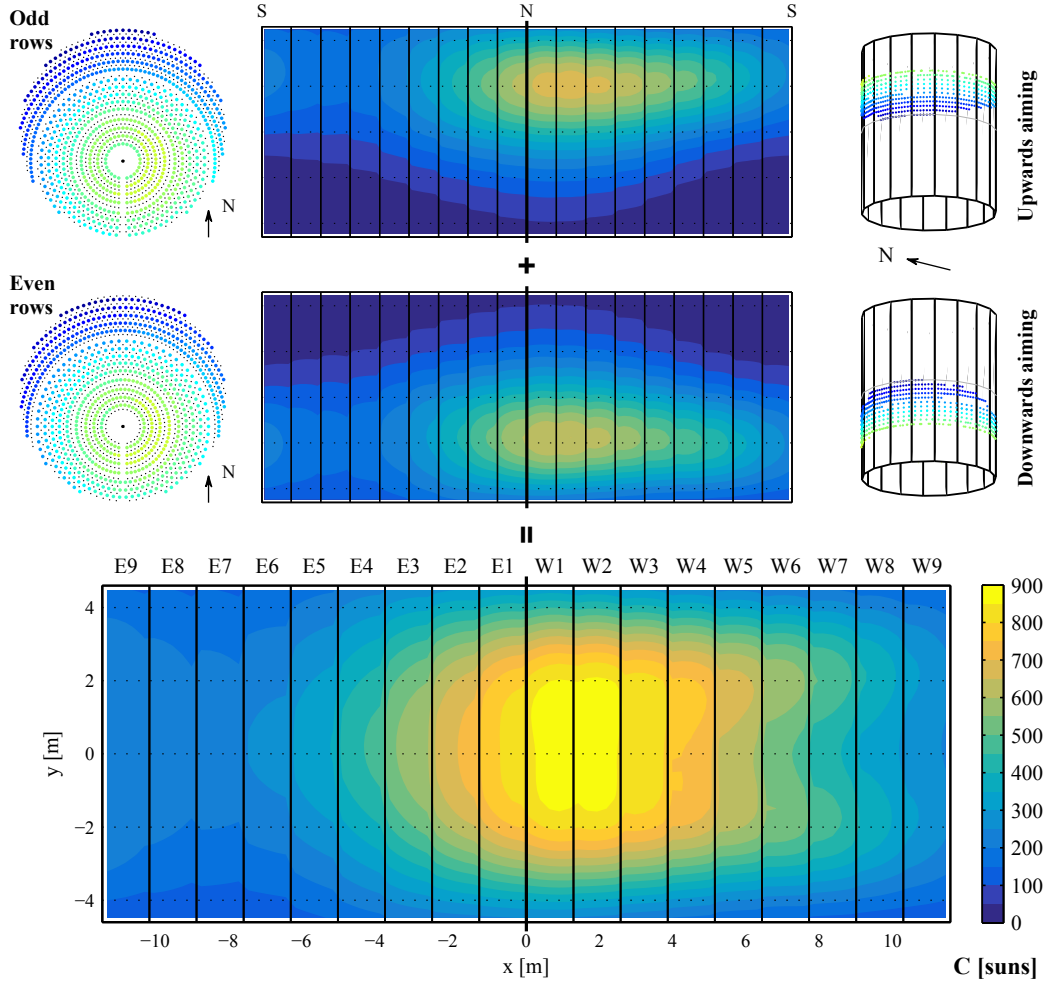


Figure 6: Symmetric aiming strategy procedure. Example for Dunhuang plant at equinox 8:00 solar time given $k = 1.5$.

251 line. Fig. 2 represented, in the 18-panels Dunhuang receiver, the 37 aim
 252 levels colored for identification in heliostat field plots of aim points.

253 In this instance, the aim levels match the receiver mesh where the flux
 254 distribution is computed. This way, or taking a vertical discretization that
 255 is submultiple of the distance between aiming levels, is advantageous for the
 256 flux map shifting procedure reported in Appendix A.

257 *5.2. Case study results*

258 For Dunhuang at equinox noon, Fig. 7 illustrates the maps of concentra-
 259 tion ratio in the west receiver panels for aiming factors equal to: 3, 2, 1, and
 260 0; keeping the same colormap scale for proper comparison. In the right side
 261 of the Figure, the mean vertical concentration profiles for each west panel is
 262 drawn according to the panel color coding at the top of the maps. In the
 263 left, the west side of the field is colored by the aim level of each heliostat,
 264 following the color levels in previous Fig. 2, which again have been slightly
 265 marked at the right side of the flux map.

266 Concentration map for $k = 3$, which is equivalent to equatorial aiming
 267 in Dunhuang case study, has a remarkable hot strip in the equatorial belt,
 268 reaching a peak concentration of 1683 suns in the most irradiated panel (W1).
 269 On the contrary, for $k = 0$ two peaks (up to 863 suns) emerge in the receiver
 270 edges, which is a very inefficient aiming with less than 50% interception; see
 271 Table 1.

272 Aiming factors between the extremes, i.e. 1 and 2, leads to flux maps also
 273 in between. For $k = 1$ in Fig. 7, the two shoulders on the vertical profiles
 274 have just appeared. Meanwhile, $k = 2$ generates a single peak, definitely
 275 smaller than with $k = 3$, of 1372 suns in panel W1. However the average
 276 concentration ratio in the whole receiver is 559, slightly smaller than the 569
 277 suns with $k = 3$.

278 From the inspection of the flux profiles, it can be noted not only the de-
 279 creasing peak alongside the panels, but also the changing pattern depending
 280 on the aiming factor. One may predict that an aiming factor between 1 and
 281 2 should lead to the flattest profiles.

282 Table 1 summarizes the spillage efficiency of Dunhuang field at equinox
 283 noon for several aiming factors ranging from 3 to 0. With symmetric aiming,
 284 η_{sp} is in between the values for unfeasible upwards and downwards aiming.
 285 Since the 68-95-99.7 rule for normal distribution was taken as the starting
 286 point for the aiming factor strategy (Subsection 2.1), the probability for
 287 Gaussian distribution in correspondence with k is displayed in the last two
 288 columns of the Table. Even though η_{sp} and Gaussian probability exhibit the
 289 same trend, it is obvious that intercept efficiency is around 50% for $k = 0$ in
 290 contrast to 0% probability for $0 \cdot \sigma$. Nonetheless, the flux distribution from
 291 the heliostats is clipped by the finite apparent shape of the receiver.

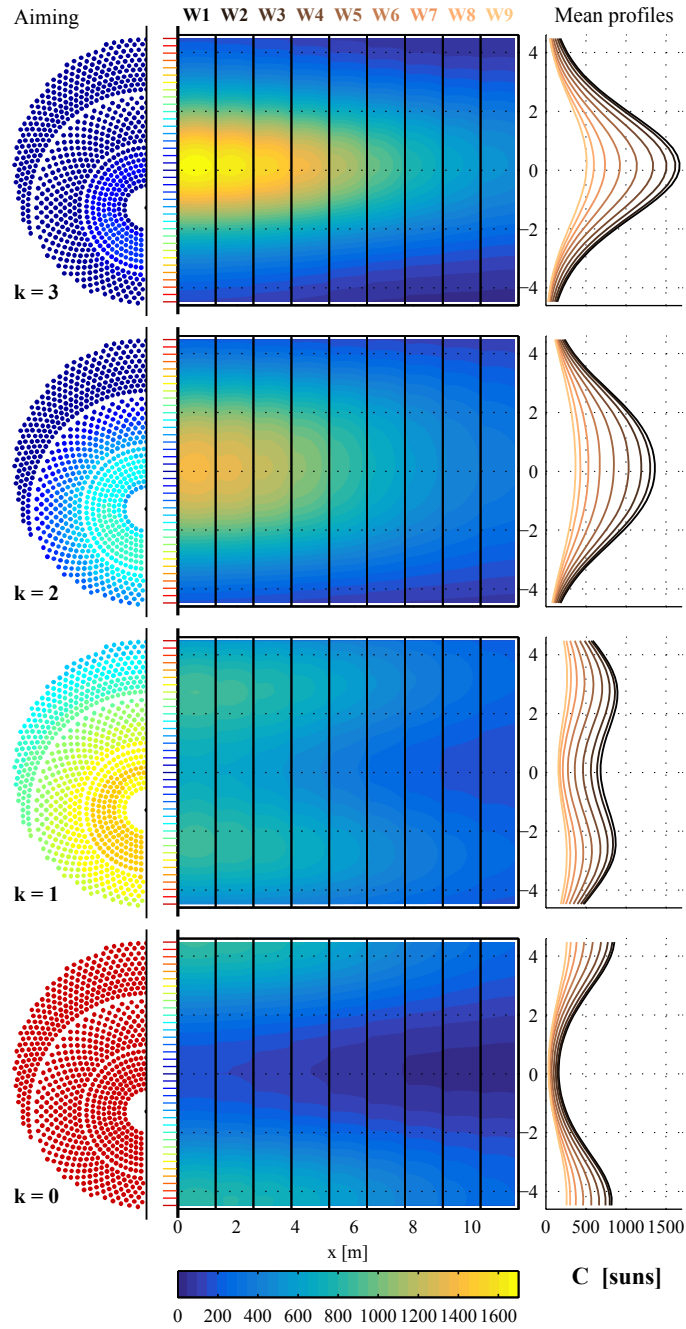


Figure 7: Maps of concentration ratio in west panels of Dunhuang receiver at equinox noon, for aiming factors equal to: 3, 2, 1, and 0. Right: Mean vertical profiles for each panel. Left: aim level for each heliostat in the field.

Table 1: Field spillage efficiency in Dunhuang at equinox noon for different aiming factors and considering: upwards, downwards, and symmetric aiming.

Aiming	η_{sp} [%]			Normal dist.	
	Upw.	Sym.	Dow.	Prob. [%]	Range
Equatorial		91.0		100.0	∞
3	90.9	91.0	91.0	99.7	$3\cdot\sigma$
2.5	90.3	90.5	90.7	98.8	$2.5\cdot\sigma$
2	88.6	89.1	89.7	95.5	$2\cdot\sigma$
k 1.5	84.0	85.2	86.5	86.6	$1.5\cdot\sigma$
1	74.6	76.8	79.0	68.3	$1\cdot\sigma$
0.5	60.5	63.4	66.4	38.3	$0.5\cdot\sigma$
0	43.7	46.8	50.1	0	0

292 6. Flat distribution

293 From the effect of different aiming factors in the symmetric aiming strat-
 294 egy, it can be concluded that there might be a k_{flat} factor where the flux
 295 distribution would be the flattest possible. In fact, given a solar position,
 296 rather than a single k_{flat} for the whole field, it can be found an specific one
 297 for each sector/panel, as described in the following.

298 6.1. Aiming by field sectors

299 The field of heliostats is divided in sectors in accordance with the target
 300 panel. Since only vertical shifting of aim points is considered, the aperture
 301 angle of the receiver panels bounds the field sectors. So that in the 18-panels
 302 Dunhuang receiver, the angle of each sector is 20° . In previous Section 3
 303 Fig. 3 represents the 18 field sectors, which are labeled according to their
 304 target panel. The heliostats in each row-sector are linked with arc lines in
 305 that Figure.

306 Distinct aiming factors for each field sector can now be defined. Then,
 307 given a k_{sector} value in a field sector, the heliostats in the same row-sector
 308 will point to the same aim level. This is explained by the almost identical
 309 beam size for the heliostats in a row-sector (Eq. 4), that, from Eq. 5 and
 310 shifting quantization, finally results in the same aim level.

311 To ensure that all the heliostats in a row-sector point to the same aim
 312 level, the beam size may be computed in two alternative ways. First, it could

313 be considered the quadratic mean of the heliostats beam radii in such row-
 314 sector. This approach was used in Ref. [17], but the correspondence between
 315 aiming factor and spillage losses, as analyzed in previous Section 4, is lost.
 316 Second, it is more appropriate to compute the beam size as the mean beam
 317 radius of the heliostats in the row-sector, as considered in the present study.

318 *6.2. Flat aiming factors: k_{flat}*

319 This Subsection describes the procedure to determine k_{flat} ; that is: the
 320 aiming factor for each field sector in order to procure the flattest distribution
 321 using the symmetric aiming strategy.

322 The purpose of finding k_{flat} is worked out by an algorithm that seeks along
 323 a sequence of k factors for those providing the flattest vertical profiles in each
 324 receiver panel. Since the flux distribution in each panel is mainly affected
 325 by the heliostats in its field sector, the search is performed sector-by-sector.
 326 And, at the end, considering the contribution of the whole heliostat field, it
 327 is checked that the vertical flux profiles are still kept flat.

328 The success of the algorithm depends on a proper selection of the sequence
 329 of k values to be swept. From the analysis in Section 4, it has been concluded
 330 that spillage losses logarithmically increase with the aiming factor. There-
 331 fore, a sequence of logarithmically spaced k factors would result in roughly
 332 constant steps in the spillage losses. On the other hand, the number of k
 333 values in the sequence should be such that nearly all aim levels were swept.
 334 It has been found that a series of around half the number of aim levels (37
 335 in this instance), performs a regular scanning. Hence, in the Dunhuang case
 336 study 18 bins have been considered. Selecting $k = 3$ (equatorial aiming) and
 337 $k = 0.5$ (almost edge aiming) as upper and lower bounds of the sequence,
 338 the whole series is: 3.00, 2.72, 2.46, 2.23, 2.01, 1.82, 1.65, 1.49, 1.35, 1.22,
 339 1.11, 1.00, 0.91, 0.82, 0.74, 0.67, 0.61, 0.55, 0.5.

340 To find out the k_{flat} of each sector, it is computed in each panel the
 341 vertical profiles resulting from each k factor following the sequence. The
 342 latest aiming factor before obtaining two peaks in the profile, is finally the
 343 k_{flat} for such sector. This algorithm is similar to the sweep routine employed
 344 in Ref. [17].

345 For the Dunhuang case study in the equinox day at 8:00 solar time, Fig. 8
 346 represents the vertical profiles of C concentration ratio for each panel in the
 347 receiver. Each grayscale profile corresponds to a k value in the sequence, as
 348 labeled in the color legend at the right side of the Figure. The red thick line
 349 in each panel identifies the flattest profile and its corresponding k_{flat} factor.

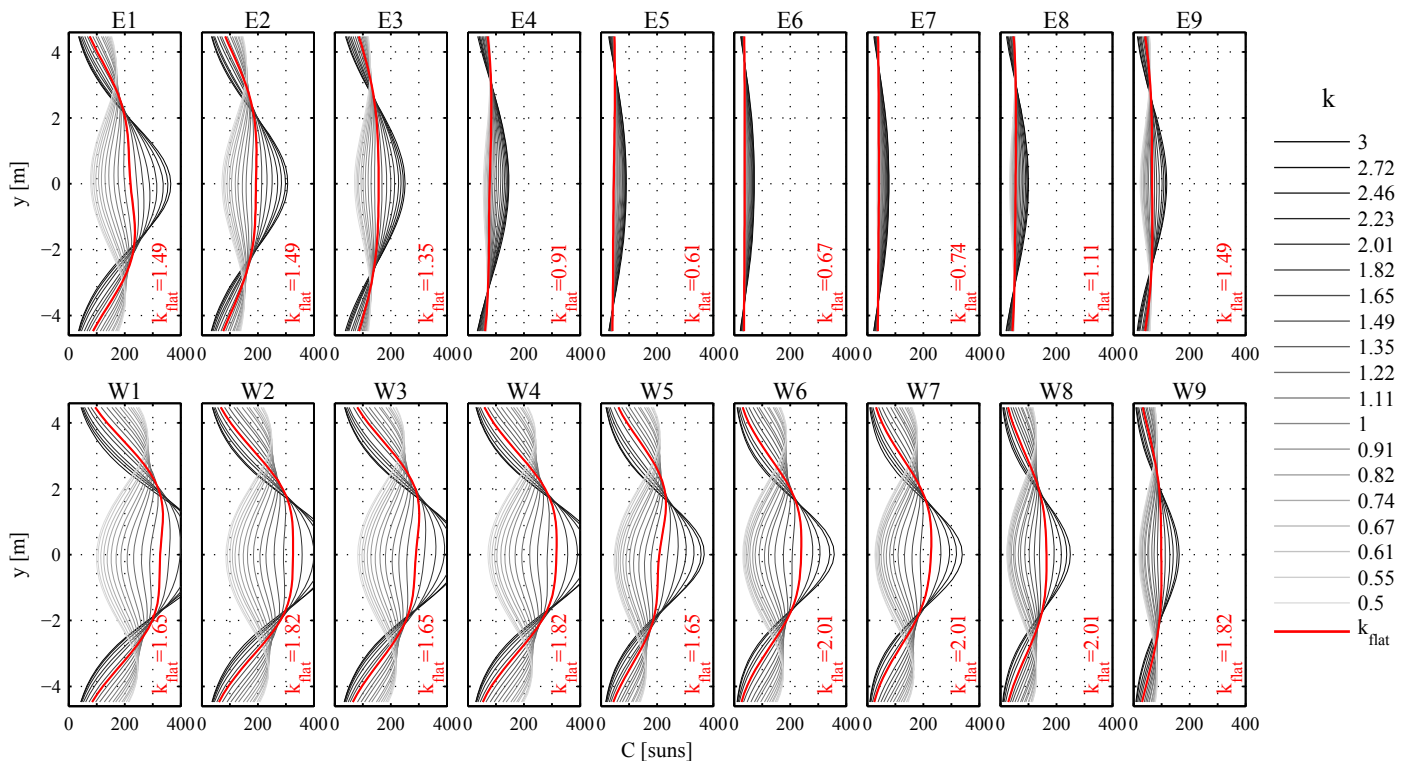
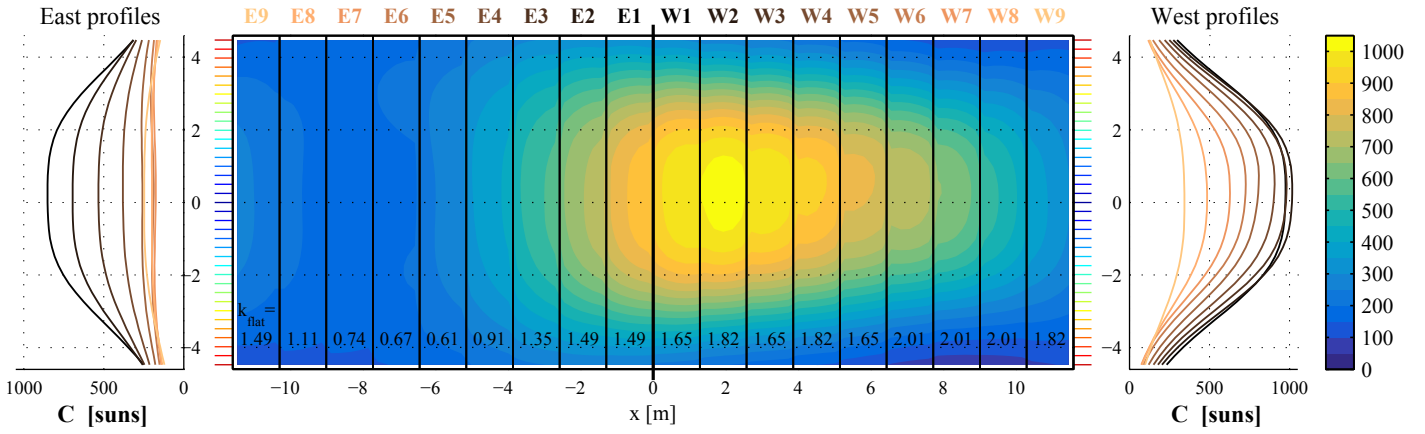


Figure 8: Vertical profiles of concentration ratio in each receiver panel of Dunhuang at equinox 8:00 solar time, for aiming factors ranging from 3 to 0.5 (sweep). Red thick curves represent the flattest profile for each panel, achievable with k_{flat} .

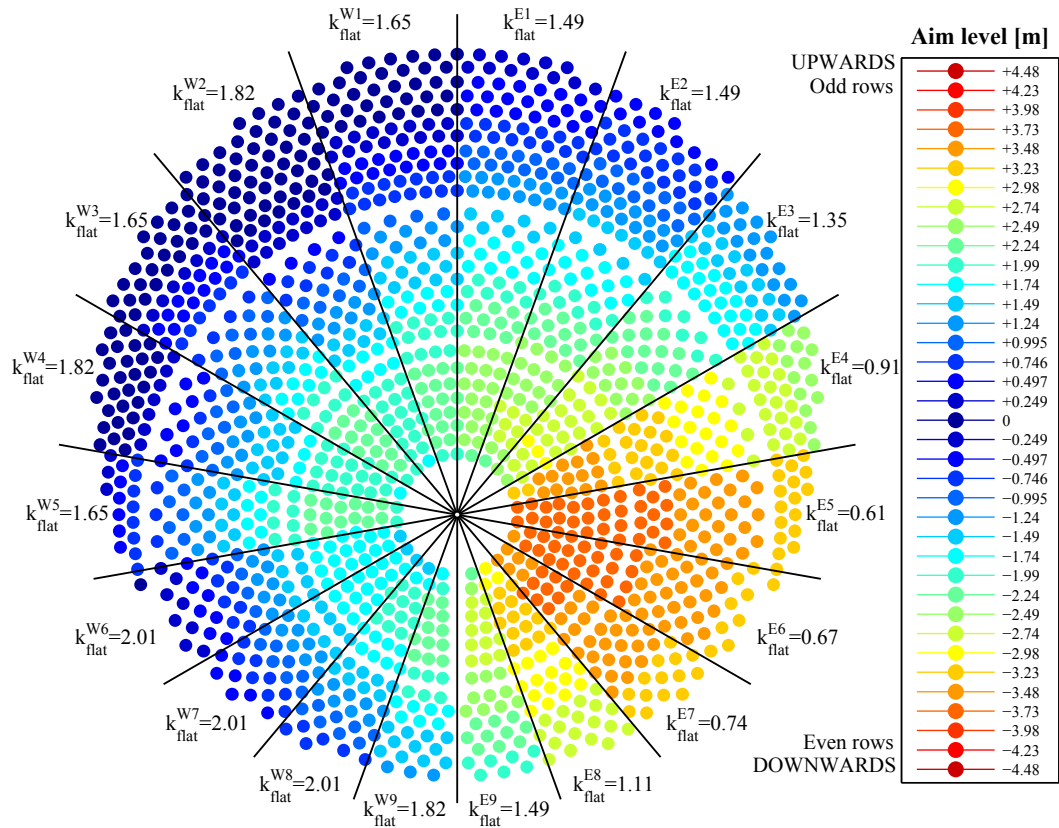
350 As can be seen in Fig. 8, the profiles for k values smaller than k_{flat} have
 351 a central valley and two peaks towards the edges, so that spillage losses
 352 gradually increase. On the contrary, aiming factors higher than k_{flat} lead to
 353 a central peak and high gradients along the receiver vertical direction.

354 The convolution-projection method used to compute the flux distribution
 355 [16] is faster than MCRT simulations. However, in computationally intensive
 356 procedures such as aiming optimization [17] or the current sweep algorithm,
 357 where diverse aiming configurations are required, significant time can be
 358 saved using the flux map shifting approach described in Appendix A. The
 359 freeware FluxSPT tool [22] also takes advantage of the mesh shifting to
 360 instantly compute the flux maps resulting from different aiming factors.

361 For the flux profiles of each receiver panel in Fig. 8, from which the k_{flat}
 362 has been obtained, just only the corresponding field sector was aimed to



(a) Map of concentration ratio and vertical profiles.



(b) Aim point level of each heliostat.

Figure 9: Dunhuang at equinox 8:00 solar time, using flat aiming factors k_{flat} .

363 the panel. By superimposing the flux distribution resulting from each sector
 364 k_{flat} , it is gained the whole *flattest* concentration map, reproduced in Fig. 9a.
 365 In the left and the right sides of the Figure, they are depicted the vertical
 366 profiles for the respectively east and west panels. From the inspection of
 367 Fig. 9a, it is noticed that the profiles resulting from the whole heliostat field
 368 are also flat in their central region; therefore, the sector-by-sector approach
 369 in the sweep algorithm to find out the k_{flat} values appears plausible.

370 On Dunhuang layout, Fig. 9b represents the aim point level for each he-
 371 liostat, that leads to the flat distribution in Fig. 9a. For the equinox 8:00
 372 case, the greatest aim point shifting about the receiver equator, identified by
 373 warm colors, takes places in the heliostats around the solar azimuthal loca-
 374 tion, specifically in sectors E5, E6 and E7. In fact, these three panels yield
 375 the lowest k_{flat} values: 0.61, 0.67 and 0.74, respectively. On the contrary,
 376 heliostats in the last rows of the anti-sun location are aimed to the equator,
 377 or close to it.

378 6.3. Case study annual results

379 Running the sweep algorithm along the year, it can be collected for any
 380 solar position the k_{flat} factor for each field sector that generates the flattest
 381 flux distribution using the symmetric aiming approach. On an azimuth and
 382 elevation diagram for Dunhuang, Fig. 10 displays for each east sector the
 383 k_{flat} using a grayscale contour map that designates the aiming factors in the
 384 sweep sequence. Red lines represent the sun path in the 21st of each month,
 385 identifying the solstices and equinoxes, and white dashed lines indicate the
 386 solar times. By means of Eq. 5 and shifting quantization of y_{shift} to the
 387 aim levels, from Fig. 10 it is ultimately gained the aim point level for all
 388 the heliostats in Dunhuang field in order to achieve the flattest possible flux
 389 distribution at any time.

390 In the northern E1 sector, the k_{flat} factor is fairly constant throughout
 391 the year. In the southern E9 sector, k_{flat} is quite uniform during a single
 392 day, but varies along the year, being highest in summer solstice and lowest in
 393 winter solstice. In the remaining sectors, it is noticeable the change of the flat
 394 aiming factor in the first hours of the day, from small k_{flat} values to higher
 395 ones around noon and in the afternoon. Although Fig. 10 only compiles the
 396 east sectors, by making the symmetrical about the noon vertical axis (i.e.
 397 zero azimuth) it is gained the k_{flat} for the analogous west panels.

398 In a similar way, Fig. 11 displays the maximum concentration ratio of flux
 399 density, C_{max} , in each east panel resulting from the application of k_{flat} aiming

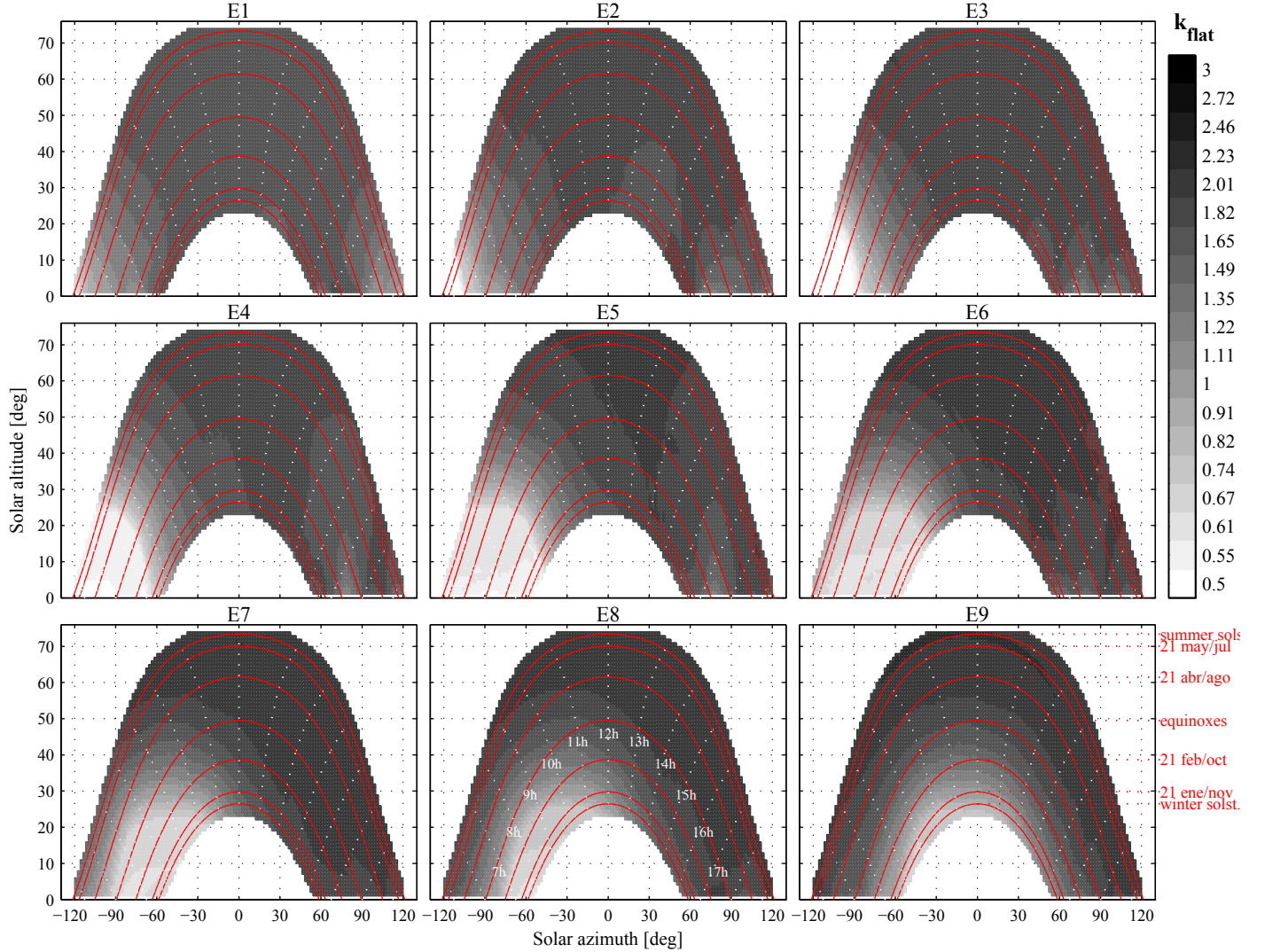


Figure 10: Flat aiming factors in Dunhuang plant for east sectors along the year (solar azimuth and elevation diagram).

400 factors. Therefore, such C_{max} is essentially the concentration ratio in the
 401 central region of the panel, where a flat vertical profile exists. As expected,
 402 the highest C_{max} values take place in the northern panels. Considering a
 403 particular panel, C_{max} is greater around noon than near sunrise and sunset.

404 Known the instantaneous DNI , from Fig. 11 and using Eq. 3 it is directly
 405 gathered the maximum flux density for each panel at any time. Whether

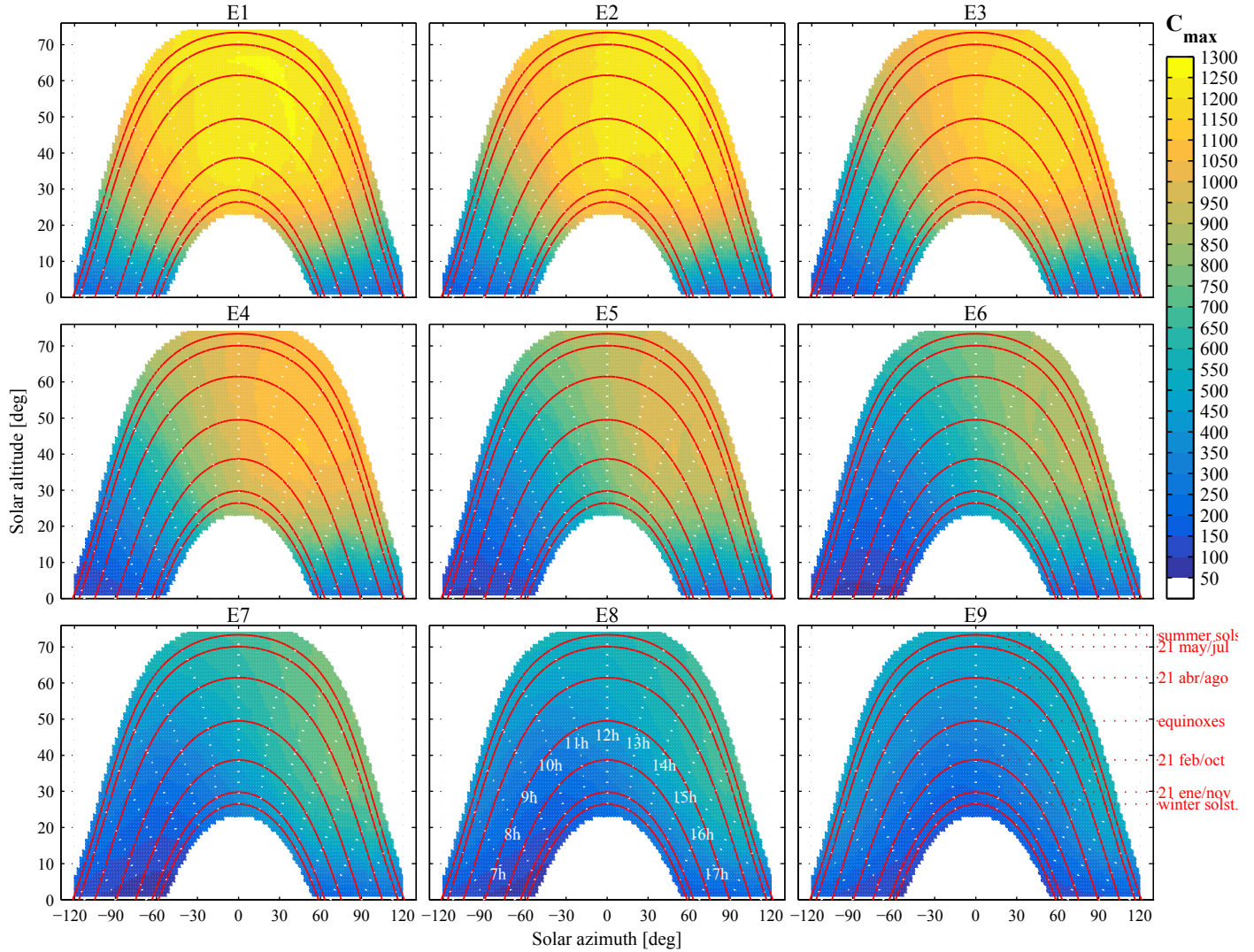


Figure 11: Maximum concentration ratios (using flat aiming factors) in Dunhuang plant for east receiver panels along the year (solar azimuth and elevation diagram).

406 limitations on the flux density exists, e.g. via Allowable Flux Density (*AFD*)
 407 [15, 17], it would be automatic to check such requirements. If, despite the
 408 flat symmetric aiming strategy, the *AFD* were exceeded, then an alternative
 409 aiming procedure, such as the fit algorithm in Ref. [17], could decrease the
 410 maximum flux at the expense of increasing spillage losses.

411 7. Conclusions

412 A symmetric aiming strategy has been developed to produce flux maps
413 symmetric about the equator in cylindrical receivers. A plethora of flux dis-
414 tributions can be generated on the basis of a single parameter: k , ranging
415 between 0 and 3 because of its analogy with Gaussian distribution probabili-
416 ty. Equatorial aiming is commonly equivalent to $k = 3$, and $k = 0$ means
417 alternatively aiming to top and bottom receiver edges. Aiming factors be-
418 tween the limits yield to symmetric maps also in between. This approach has
419 been implemented in the software tool FluxSPT, available for free download
420 from [22].

421 The k aiming factor serves to estimate the size of the beam reflected by
422 each heliostat in a field, and then vertically shift to the target aim level
423 so that the spots are alternatively tangent to either the upper or the lower
424 edge of the receiver. It has been found that interception is logarithmically
425 connected to k factor.

426 To reduce flux gradients and thermal stresses along the receiver tubes,
427 a flat vertical profile is convenient. The k_{flat} values leading to the flattest
428 possible symmetric flux maps have been deterministically found using a sweep
429 methodology. This search becomes faster with the accurate enough mesh
430 shifting procedure in the Appendix.

431 Taking advantage of sectorial division of the heliostat field in correspon-
432 dence with the receiver panels, specific k_{flat} values have been provided for
433 each sector, so that the flux profile is made flat in every single panel. For
434 Dunhuang SPT plant in China, flat aiming factors for the whole year have
435 been summarized in azimuth and elevation diagrams; k_{flat} is quite constant
436 along the year except near sunrise and sunset, respectively in east and west
437 field sectors.

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442 Appendix A. Flux map shifting

443 The computation of the flux maps for each aiming factor, which is per-
444 formed with the convolution-projection method in Ref. [16], in principle re-

445 quires an individual calculation for every aim point level considered. This
 446 Appendix puts forward a methodology to quickly determine the flux distri-
 447 butions by an heliostat pointing to any aim level in the receiver, given a
 448 specific solar time. This procedure, which saves considerable computation
 449 time in sweep and optimization routines, involves a negligible error com-
 450 pared to direct, long-lasting, computation; as also justified at the end of this
 451 Appendix.

452 The proposed methodology consists in computing only once the flux dis-
 453 tribution by a heliostat, as in Ref. [16], considering equatorial aiming. Such
 454 flux map is saved and, when a different aim level for the heliostat is required,
 455 the original flux map is vertically shifted a distance equal to that between the
 456 equator and the target level. Fig. 12 graphically illustrates this procedure,
 457 which is implemented in the FluxSPT software tool [22].

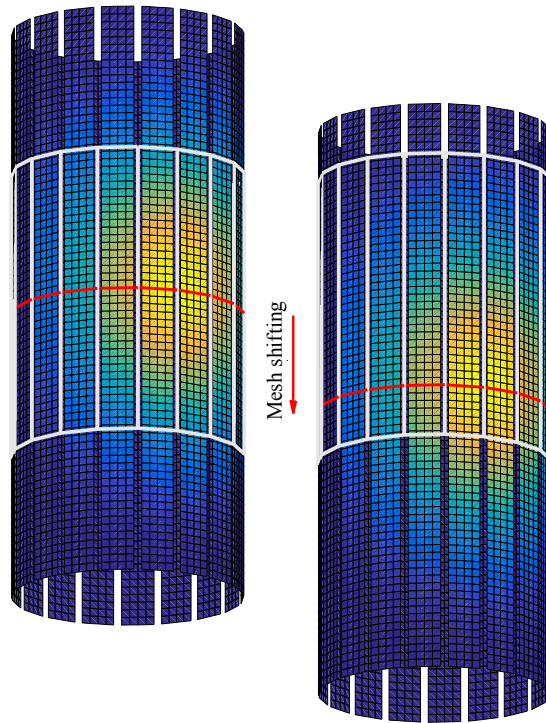


Figure 12: Graphical explanation of the flux map shifting procedure. Left: mapping on extended receiver mesh for equatorial aiming. Right: mapping shifted downwards for a given k -factor of 0.5. (Example at equinox noon for heliostat in last row, north azimuth).

458 Operationally, the flux map shifting procedure entails two issues to con-

459 sider. Firstly, the computational mesh in the receiver has to be extended
460 beyond its upper and lower borders. Since the maximum shift would be half
461 the receiver height ($k = 0$, edge aiming), the height of the full mesh might
462 be twice the receiver height. This is shown in Fig. 12, where the receiver is
463 highlighted by white lines.

464 Secondly, the computational mesh should match the predefined aim levels
465 so that interpolation is favorably avoided. This can be achieved by taking a
466 nodal spacing in the receiver mesh equal to, or submultiple of, the distance
467 between aim levels. For instance, the 37 aim levels match the vertical dis-
468 cretization in Fig. 12, then the flux map shifts from the equator (red line in
469 left image) to the target aim level (red line in right image).

470 Comparing the direct computation with the proposed flux map shifting
471 procedure, very slight differences are found. For the extreme case of $k = 0$,
472 Fig. 13 points out, for each heliostat in Dunhuang field at equinox noon, the
473 associated error in the calculation of flux maps and spillage. The relative
474 error in the spillage loss factor (east field side) is inversely proportional to
475 the slant range, so that for far heliostats the relative error is around zero
476 and for near heliostats almost 3%. In the west side of Dunhuang layout in
477 Fig. 13, it is represented the root mean square deviation in the concentration
478 maps, $RMSD(C)$, which is maximum in the near northern heliostats. This
479 fact is explained by the highly concentrated spot resulting from the heliostats
480 closest to the tower.

481 For the heliostat in the first row and north azimuthal position, Fig. 14
482 displays the concentration maps using both direct computation and flux map
483 shifting, respectively in the eastern and western receiver panels. Even though
484 the highest $RMSD(C)$ at noon takes place with that heliostat, the differences
485 between both computation methods are insignificant; $RMSD(C) = 0.0445$
486 suns. It is concluded that computational cost can be saved with the flux map
487 shifting procedure, which may be utilized in aiming optimization without
488 meaningful loss of accuracy.

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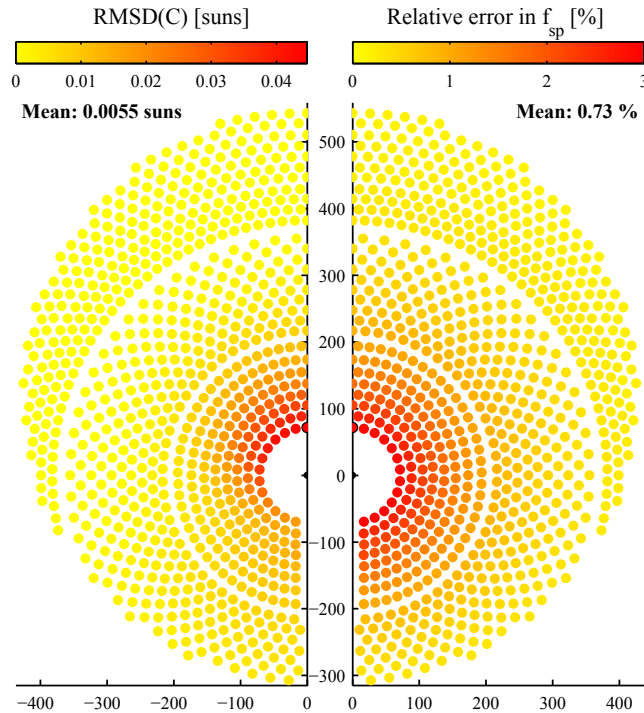


Figure 13: Errors because of flux map shifting respect to direct computation. Root mean square deviation in concentration map (west half) and relative error of spillage loss factor (east half) in Dunhuang field at equinox noon, given $k = 0$ (downwards aiming).

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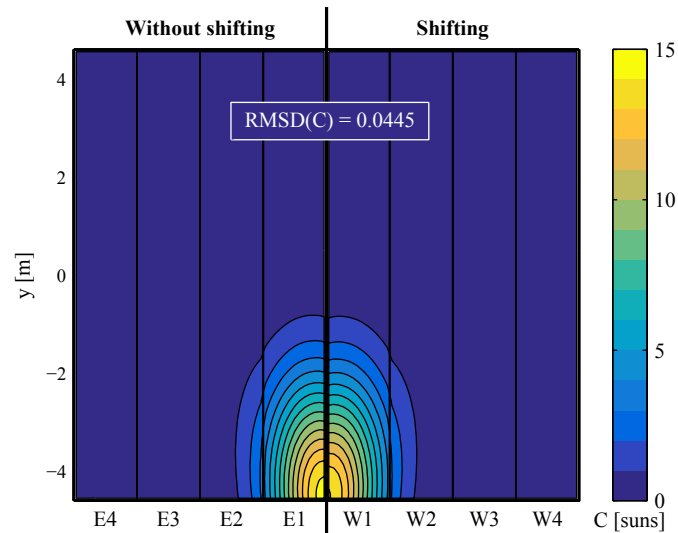


Figure 14: Concentration map at equinox noon for a heliostat (first row, north azimuth), given $k = 0$. Left/East side of the receiver: direct computation. Right/West side of the receiver: mesh shifting.

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